The LHCb detector

Int.J.Mod.Phys. A 30, 1530022 (2015)

- Fully instrumented in forward region $2 < \eta < 5$
- Excellent vertex resolution
 - $B_{\rm s}$ oscillation at 40 fs (average boost $\beta \gamma \sim 3$)
- Excellent mass resolution
 - 0.5% at $m(\Upsilon \rightarrow \mu \mu)$
- Cherenkov PID capabilities
- Good jet reconstruction
 - 10-20% energy resolution for jets with $p_{\rm T}$ >10 GeV
 - *b*(*c*) tagging eff 65%(25%) for 0.3% contamination





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The LHCb detector

- Lower luminosity (and low pile-up)
 - 1/8 of ATLAS/CMS in Run 1
- Capable of very soft triggers!
 - At hardware level (L0):
 - ▶ $\epsilon = 95\%$ for detached $\mu\mu$ with $p_{\rm T} > 1 \text{GeV}/c$
 - Calo trigger at ~3.5 (~2.5) GeV for hadrons (electrons)
 - At Software level (HLT):
 - Topological triggers on detached vertices
 - Also PID and jets in trigger!
- Trigger-less upgrade (2020)
 - Read-out detector in real time
 - Will trigger on detached vertices at first level!



Direct searches at LHCb

- **Complement** ATLAS/ CMS searches in certain phase space regions :
 - Light masses
 - soft trigger and forward acceptance
 - Low lifetimes down to 1 ps
 - excellent vertexing and boost

Increasing interest in direct searches!

- 1. Produced in *B*/*D* decays (prompt / long-lived)
- 2. Produced in *pp* collision (prompt / long-lived)

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FIG. 3 in refe rana n muons also sh althoug the cou the cou Univer

FIG. 2. Diagrams for the indirect (top) or direct (bottom) approach in searches. In the top diagram, a Majorana neutrino is produced off-shell in a $D_{(s)}^+$ decay to a final state with two same-sign muons (with the same diagram, the Majorana neutrino could be also produced on-shell). In the bottom one, a hidden valley pion is produced on-shell to later decay to a

LHO 3jorana of the

LLP at LHCb



<u>Tracks from long-lived in LHCb:</u>

- Within VELO (< 50 cm)
 - in reality more like < 20 cm
- **Up to TT** (<200 cm)
 - Worse vertex and *p* resolution ($K_{\rm S}(\pi\pi)$ resolution 2× larger)
 - Not available in trigger (studies ongoing)

- VELO envelope at ~5 mm from beam
 - Detailed material veto is used
 - <5 mm: background mainly from heavy-flavour background
 - >5 mm: background mainly from material interaction

Direct searches in heavy flavour decays

Hidden Sector in $B \rightarrow K^{(*)}\chi(\mu)$

Phys Rev Lett 115 161802 (2015) Phys Rev D 95, 071101(R) (2017)

- Look for new hidden-sector bosons in $b \rightarrow s$ penguin transitions
 - Can be axion or (long-lived) inflaton
- LHCb collected world record samples of rare decay $B \rightarrow K^{(*)} \mu \mu$
- Allow detached $\mu\mu$
- MVA selection independent of $m(\mu\mu)$ and τ (uBDT)
- Search for narrow peak in $m(\mu\mu)$, excluding QCD resonances



Hidden Sector in $B \rightarrow K^{(*)}$

10³

- BR normalised to rare SM decay
- Constraint set on lifetimes [0.1-1000] ps (30µm to 30cm)
- $\tau(\chi)$ [ps] LHCb Run 1 95% CL 10² $B^+ \rightarrow K^+ \mu \mu$ 10 10⁻¹ 1000 2000 3000 4000 $m(\gamma)$ [MeV/c²] 10^{-4} Theory LHCb 95% CL 10⁻⁵ $\mathbb{Z} B^+ \rightarrow K^+ \gamma$ $\theta^2 \left[rad^2 \right]$ 10^{-7} Theory 10⁻⁸ **Cosmological constraints** 2×10⁻¹ 2 3 4 5 m_{γ} [GeV/c²]

Phys Rev Lett 115 161802 (2015)

 $ightarrow \mathsf{K}^+ \chi) \times \mathsf{BR}(\chi$

BR(B⁺

7

Phys Rev D 95, 071101(R) (2017)

- Constrain on new scalar mixing $|\text{Higgs}\rangle + \cos\theta |\text{Higgs}\rangle$ $|\chi\rangle_{\rm phys} = \cos\theta |\chi\rangle + \sin\theta |{\rm Higgs}\rangle$ $\tau \propto 1/\theta^2 \qquad \mathcal{B}(B^+ \to K^+ \chi) \propto \theta^2$
- Nearly rule out the inflaton parameter space below $2 m_{\tau}$

Majorana neutrinos in $B^- \rightarrow \pi^+ \mu^- \mu^-$

Phys Rev Lett 112 131802 (2014)

- Lepton number violating $B^- \rightarrow \pi^+ \mu^- \mu^-$ can proceed via on-shell Majorana neutrinos
- Look for *B* mass peak, then extract limit as a function of m_N
- Limit set on N($\pi\mu$) lifetimes up to 1000 ps
- Constraints on mixing angle $V_{\mu 4}$
 - Recently revisited B Shuve, ME Peskin, Phys.Rev. D94 (2016) no.11, 113007
- Searches in other B/D channels foreseen
- Can also search using W \rightarrow jet $\mu^{-}\mu^{-}$





Produced in pp collisions



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$LLP \rightarrow \mu + jets$

Eur. Phys. J. C (2017) 77:224

Result interpreted in various models:



candidate from the safety dataset. The two tops selected from the background region of the muon

$LLP \rightarrow \mu + jets$

- Re-interpreting in terms of Sterile Neutrinos
- Recasting search with simplified implementation of the analysis
- Limit is promising, but has to be taken with a pinch of salt
- Future searches:
 - can add prompt lepton for triggering
 - can look into e and τ as well





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$\mathrm{LLP} \to \mathrm{jet} \,\, \mathrm{jet}$

LHCb-PAPER-2016-065 very soon in ArXiv

- Tested the region: m_{π} =[25-50] GeV, τ =[2-500] ps
- Example: for $m_{\pi} = 50$ GeV exclude BR > 30% for $\tau = [5-50]$ ps ($c\tau = [1.5-15]$ mm)
- Contributing at low lifetime and low mass despite lower luminosity and acceptance!





LLP to jetjet searches

Areas where $H \rightarrow \pi_V \pi_V$ BR is smaller than 50% (at 95% CL)



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LLP at upgraded LHCb

- Removing L0 !!
 - Can trigger directly on displaced vertices (or even more complex signatures)
 - Also jets are in the trigger
 - No $p_{\rm T}$ requirements —> low masses
- Upgraded Vertex Locator
 - much faster —> vertices in the trigger
 - more precise —> less ghost tracks
 - thinner box —> less material interaction
- Will concentrate efforts in low mass, low lifetimes



Low mass di-lepton program

Dark Photon A' (auto-normalising to γ^*)

Low-mass CP-odd Higgs (*a*)



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Other low-*µ*

The scale of m_{ν} is not measured directly, as net trino oscillation experiments probe only the squared mas splittings, Δm_{ν}^2 . The actual values of m_{ν} can vary froi massless (which is a viable option only for the lighter mass eigenstate) to the upper bounds supplied by comology ($m_{\nu} \leq 0.23 \text{ eV}$) [12] and direct neutrino masearches. ($m_{\nu} \leq 2 \text{ eV}$) [13]. For the heavier mass eigen states, a lower bound is given by the experimentally determined squared mass splittings. For both the norm: and inverted hierarchy at least one mass eigenstate musbe heavier than $\sqrt{\Delta (m^2)^{\text{atm}}} \simeq 0.05 \text{ eV}$, giving a lower bound on the mixing angle. From the see-saw relation i Eq. (4), the expected value of the mixing angle is: q

• Pair of \mathfrak{P}_{s} ghts has $1^{-11} \times \left(\frac{1 \text{ GeV}}{M_N}\right)$. (5) (low mass) Bat

of m_{ν} is not on experiments a well-motivated target for experimenphasized, however, that more complicated mass genera m_{ν}^2 . The actual values of m_{ν} can vary from the lightes

ate) to the upper **Boundsass** pflitble beavy, sterile state M_N is essentially $\lesssim 0.23 \text{ eV}$ [12] and discriminateriof the model. Of particular interest to $e_e \lesssim 2 \text{ eV}$ [13]. For the heavier that eigerkinematically accessible to current bound is given by the perpendicular $M_N \ll \text{TeV}$; the RH neutrino can be ared mass splitting ctly producted in SM interactions, but the production hierarchy at least one mass fixer $M_N \ll \text{TeV}$; the RH neutrino can be an $\sqrt{\Delta(m_\mu^2)^{\text{atm}}} \approx 0.05 \text{ eV}$ Heiving a low for produced in SM interactions.



FIG. 2: *(Left):* Right-handed neutrinos (N) decay via the electroweak interactions due to mixing with LH neutrine they also decay to the Higgs via Yukawa couplings (masses). *(Right):* At low masses, $M_N \leq \text{GeV}$, the excluse hadronic decays of N, such as $N \to \pi^{\pm} \mu^{\mp}$, are relevant.



Conclusions

- Limited acceptance, smaller luminosity, no MET
- LHCb demonstrated to be extremely good at:
 - Looking for on-shell new physics from B/D decays
 - Looking for long-lived particles with low mass and short lifetime
- Big potential in low mass di-leptons: light Higgs, dark photons, ...
- Bright future ahead:
 - 3/fb in Run 1, expect 5/fb in Run 2 (low pile-up)
 - A lot of potential in "trigger-less" upgrade wertex detector tracking RICH, move to purely software trigger



BACKUP

The LHCb detector

Int.J.Mod.Phys. A 30, 1530022 (2015)



Charged Massive Stable Particles

- Select pair of muon-like tracks in mass range [120, 300] GeV/c^2
- Train Neural Network to combine RICH information with dE/dx from VELO and calorimeters
- Limit is not competitive with D0 (low mass) and ATLAS (high mass)
- Proof of concept for future searches!
- Possibly move to
 - single CMSP signature (3 times eff)
 - and/or to lower masses (??)



Long-Lived Particles at LHCb

20

 10^{-2}

 10^{-1}

 B^0

• Coverage complementary to ATLAS and CMS:

- Shorter lifetimes (down to ~1 ps)
- Lighter masses (down to ~25 GeV for di-jets)
- Bonus for LHCb:
 - can look for LLP from (rare) decays of heavy-flavour
 - can use the RICH to identify charged massive stable particles



D0

10

10²

 V_{ts}

 $m_{\rm H} = 90-200$

1

 V_{tb}

Regions with (subset of) published



Pieter David PhD thesis

10³

 K^{*0}

cτ (cm)

Charged Massive Stable Particles

EPJC 75 (2015) 595

- Charged Massive Stable Particles
 stable = can pass through the µ-stations
 Model considered:
 CUSY stars are by NLCD in mCMCP
- SUSY stau can be NLSP in mGMSB
- long-lived with m>100 GeV/c2 S Dimopoulos et al [NPB488(1997)39] GF Giudice and R Rattazzi [Phys.Rep. 332(2011)419]
- CMSP can leave a signature as:
 - Smaller energy loss dE/dx
 - Longer Time of Flight
 - Absence of Cherenkov signal
- Several experiments searched for them
 - LEP, Tevatron, HERA, ATLAS/CMS



Charged Massive Stable Particles

EPJC 75 (2015) 595

- Select pair of muon-like tracks in mass range [120, 300] GeV/c^2
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Future: Emerging Jets

- "Emerging jets":
 - Jets with many displaced vertices are smoking gun for dark parton 'shower' (models with composite dark sector)

Schwaller, Stolarski, Weiler, [arXiv:1502.05409]

- LHCb has potential
 - precise jet vertexing
 - sensitive to low mediator mass Displaced Di-Jet



