

New Physics?

The hierarchy problem of the electroweak Standard Model revisited

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Abstract

A careful renormalization group analysis of the electroweak Standard Model reveals that there is no hierarchy problem in the SM. In the broken phase a light Higgs turns out to be natural as it is self-protected and self-tuned by the Higgs mechanism. It means that the scalar Higgs needs not be protected by any extra symmetry, specifically super symmetry, in order not to be much heavier than the other SM particles which are protected by gauge- or chiral-symmetry. Thus the existence of quadratic cutoff effects in the SM cannot motivate the need for a super symmetric extensions of the SM, but in contrast plays an important role in triggering the electroweak phase transition and in shaping the Higgs potential in the early universe to drive inflation as supported by observation.



ould this be it?

New Physics?

Natural Tuning: Towards A Proof of Concept

The hierarchy proble Mo

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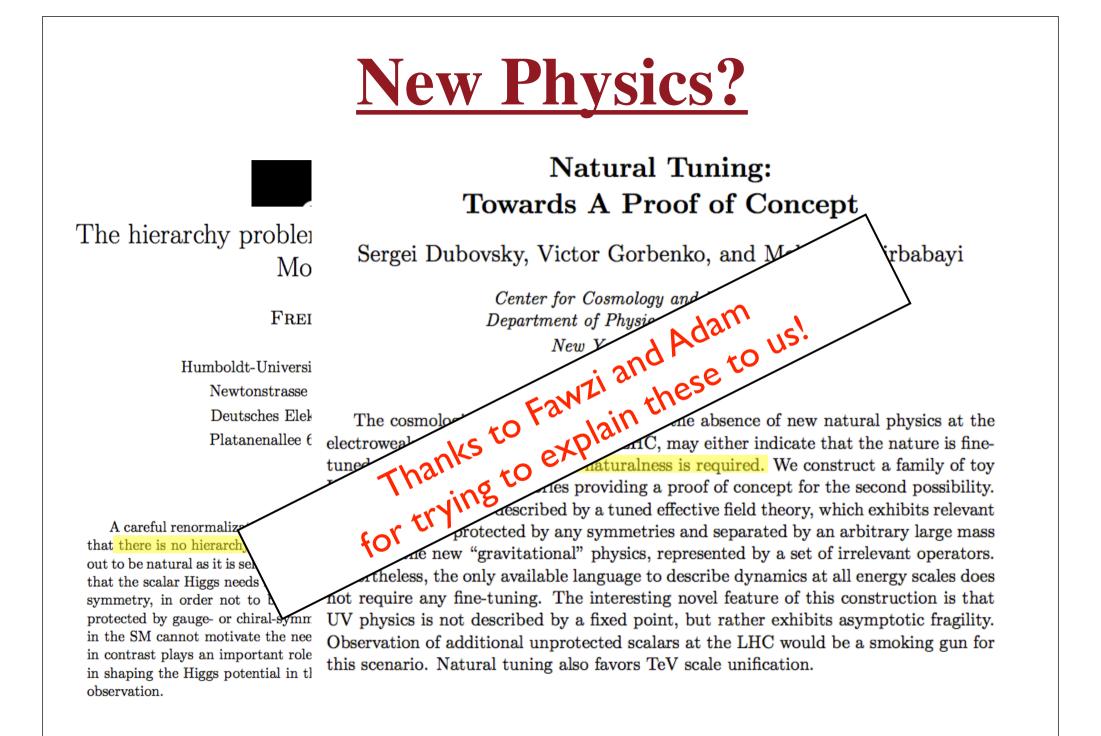
A careful renormalization group that there is no hierarchy problem out to be natural as it is self-protected that the scalar Higgs needs not be 1 symmetry, in order not to be muc protected by gauge- or chiral-symm in the SM cannot motivate the nee in contrast plays an important role in shaping the Higgs potential in the observation.

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Abstract

The cosmological constant problem and the absence of new natural physics at the electroweak scale, if confirmed by the LHC, may either indicate that the nature is finetuned or that a refined notion of naturalness is required. We construct a family of toy UV complete quantum theories providing a proof of concept for the second possibility. Low energy physics is described by a tuned effective field theory, which exhibits relevant interactions not protected by any symmetries and separated by an arbitrary large mass gap from the new "gravitational" physics, represented by a set of irrelevant operators. Nevertheless, the only available language to describe dynamics at all energy scales does not require any fine-tuning. The interesting novel feature of this construction is that UV physics is not described by a fixed point, but rather exhibits asymptotic fragility. Observation of additional unprotected scalars at the LHC would be a smoking gun for this scenario. Natural tuning also favors TeV scale unification.



Nevertheless

- Clear structure in fermionic sector unexplained
 - Evidence of some selective principle (why are there no neutral colored fermions?)
 - Proton stability, running of couplings suggestive of at least one other scale relevant to SM particles, ~10¹⁵ GeV
 - Either fine-tuning, or a closer scale

Qua	rks		
U	C	†	ags fie
d	S strange	bottom	Higgs 15 Th
	ins fie	H.	Higgs boson
electron	H	T	ages
Ve electron neutrino	V _µ muon neutrino	V _T	Hings new C
Lept	tons		

Simplified models for top partners/Vector-like quarks

 Well motivated by naturalness argument (same-spin partners of the top)

Considerable experimental effort

 Difficult to combine/interpret results (experiments are moving in this direction)

Many people involved here in Les Houches!

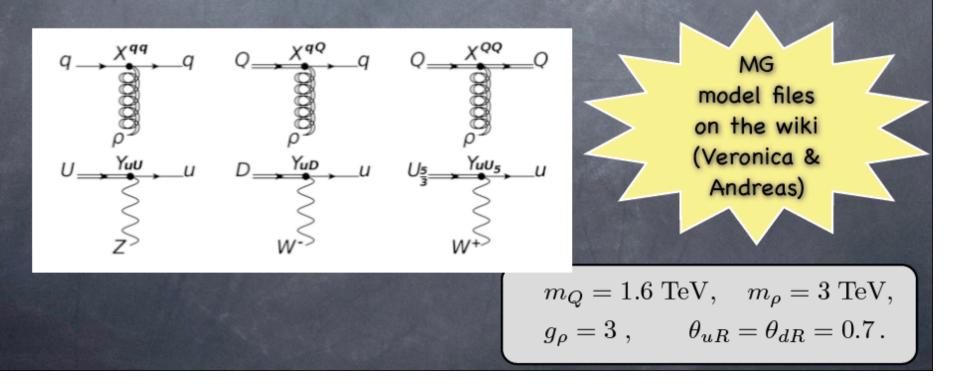
Two classes of simplified models: 'Generic' Q_V - beyond 3rd gen

- Consider a single Q_V in a given representation of SU(2)xU(1).
- Allow for couplings to light (1st) generations still in the naturalness business!
- Very limited number of parameters: m_Q, g* (coupling strength), R_L (ratio of light/3rd generation decays.
- Implementation of models in FeynRules/MadGraph/CalcHEP available on the WIKI (by G.Cacciapaglia, A.Deandrea, A.Datta)

Two classes of simplified models: MFV right-handed compositeness

M.Redi, V.Sanz, M.de Vries, A.Weiler: 1305:3818

- Model of MFV composite quarks, mainly coupling to a heavy vector and light generation
- 3-body decay into 3 jets, or 2 jets via chromomagnetic operator.

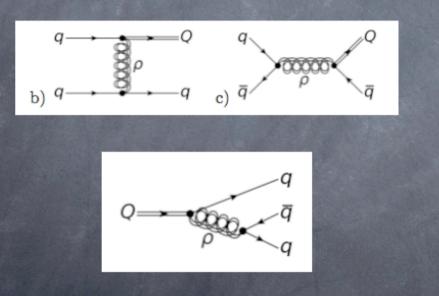


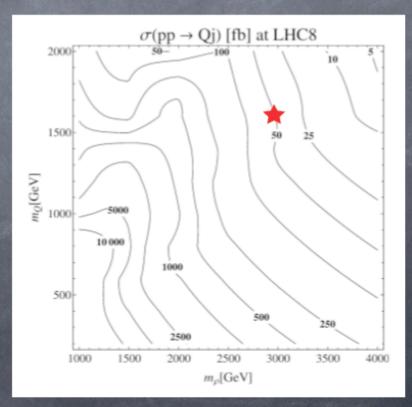
MFV right-handed compositeness

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Study of single-production in 3 jets.

V.Sanz, J.Ruiz Alvarez



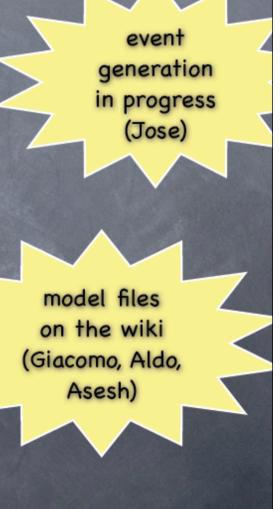


2 benchmark simplified models:

 $(1, 2/3) = T \Rightarrow T \rightarrow W^+ b, W^+ j, Zt, Zj, ht, hj$ $(2, 1/6) = \{T, B\} \Rightarrow \begin{cases} T \rightarrow Zt, Zj, ht, hj \\ B \rightarrow W^- t, W^- j \end{cases}$ $(2, 7/6) = \{X, T\} \Rightarrow \begin{cases} X \rightarrow W^+ t, W^+ j \\ T \rightarrow Zt, Zj, ht, hj \end{cases}$

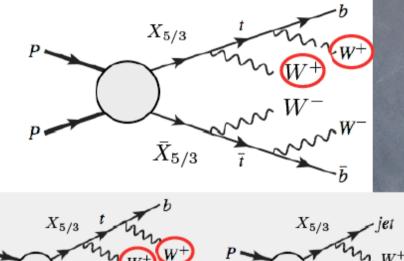
- Present in motivated models.
- Milder bounds on couplings from Flavour.

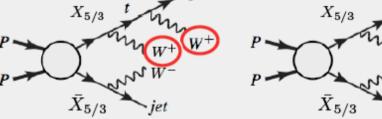
$$m_Q = 1 \text{ TeV}, \quad g^* = 0.1, \quad R_L = 0.5$$



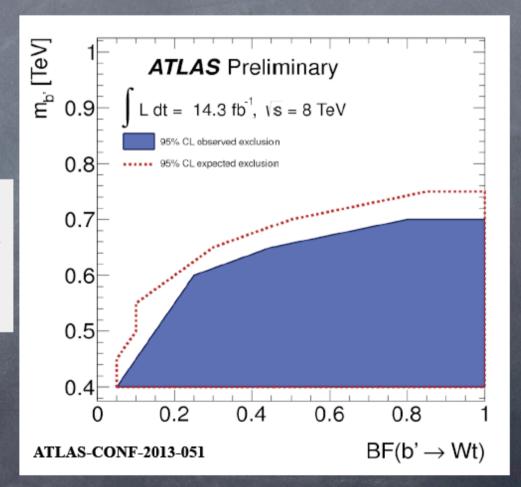
Combination of searches in Pair Production

A.Belyaev, G.Cacciapaglia, A.Datta, A.Deandrea, T.Golling + L.Panizzi





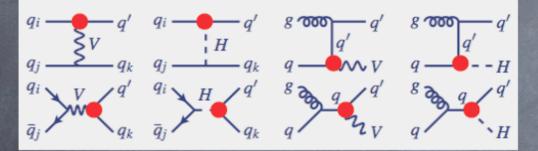
Systematic study for all cases.



Combination of searches in Single Production

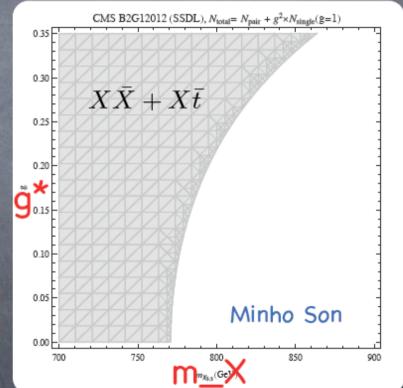
New single production channels

Single production: $pp \rightarrow q' + \{q, V, H\}$



Systematic study for all cases.

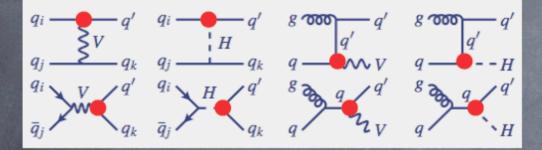
Inclusion of couplings to light generations. L.Basso, A.Belyaev, G.Cacciapaglia, A.Datta T.Golling, S.Lee, M. Son + M.Buchkremer



Combination of searches in Single Production

New single production channels

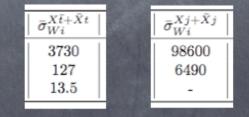
Single production: $pp \rightarrow q' + \{q, V, H\}$



Coupling to 1st generation can give large contributions to single-production:

$$\sigma(X\bar{t} + \bar{X}t) = \kappa_X^2 \left(\xi_W \sum_{i=1}^2 \zeta_i \left(\bar{\sigma}_i^{X\bar{t}} + \bar{\sigma}_i^{\bar{X}t}\right)\right)$$

1211.5663, 1305:4172



(in fb)

L.Basso, A.Belyaev, G.Cacciapaglia, A.Datta T.Golling, S.Lee, M. Son + M.Buchkremer